# Explicit Cascade Decompositions of Unitary Systems 

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Abstract. We consider the unitary linear dynamic system $\alpha=(X, U, V, A, B, C, D)$ of the form

$$
\begin{aligned}
& \frac{d x}{d t}=A x(t)+B u(t) \\
& v(t)=C x(t)+D u(t)
\end{aligned}
$$

where $x(t) \in X, u(t) \in U, v(t) \in V$. The operators $A, B, C, D$ are linear bounded and the operator

$$
T=\left[\begin{array}{ll}
A & B \\
C & D
\end{array}\right]: X \oplus U \rightarrow X \oplus V
$$

is unitary. The purpose of this paper is to decompose explicitly a unitary system by factorization of its transfer function and by its invariant subspace.

## 1. Introduction

1.1. Livsis originated and has developed the theory of linear dynamic systems in infinite-dimensional spaces and the theory of unitary, dissipative systems [10, 11]. In different mathematical languages, this theory has been studied by Nagy-Foias [12], Brodskii [6], Arov [1, 2], De Branges [5], and Khanh [8, 9].

The problem of cascade coupling or cascade decomposition of systems was posed naturally. In the case of infinite-dimensional systems, the problem was studied by Bart, Gohberg, Kaashoek [4], Wang, Davison [13], Chen Chi Tong, Doeser [7], and several mathematicians. On the other hand, Livsis [11], Arov [1, 2], Khanh [8, 9], Ball and Kriete [3], and De Branges [5], have studied this problem for unitary or passive systems.

It is indicated in $[6,12]$ that there is a one-to-one correspondence between the existence of an invariant subspace of the main operator and the regular factorization of the transfer function. Then, the problem of decomposing a unitary system by factorization of its transfer function and by its invariant subspace is determined
$[6,10,11,12]$. However, an explicit construction of these cascade decompositions is an interesting subject to consider. This is the main purpose of the paper.
1.2. Let $X, U, V$ be separable Hilbert spaces. Consider a linear discrete stationary dynamic system $\alpha=(X, U, V, A, B, C, D)$ of the form:

$$
\begin{aligned}
x_{n+1} & =A x_{n}+B u_{n} \\
v_{n} & =C x_{n}+D u_{n}
\end{aligned}
$$

where $x_{n} \in X, u_{n} \in U, v_{n} \in V$. The operators $A: X \rightarrow X, B: U \rightarrow X, C: X \rightarrow V$, $D: U \rightarrow V$ are linear bounded.

The spaces $X, U, V$ are called the state space, the input space, and the output space, respectively.

The operator function of the complex variable

$$
\theta(z)=D+z C(I-z A)^{-1} B
$$

is called the transfer function of system.
The subspaces $X_{\alpha}^{C}=\bigcup_{0}^{\infty} A^{k} B U, X_{\alpha}^{0}=\bigcup_{0}^{\infty} A^{* k} C^{*} V$ stand for controllable and observable subspaces of $\alpha$, respectively. The system is said to be controllable if $X=X_{\alpha}^{C}$, observable if $X=X_{\alpha}^{0}$, and simple if $X=X_{\alpha}^{C} \cup X_{\alpha}^{0}$.

Definition 1. Let $\alpha_{k}=\left(X_{k}, U, V, A_{k}, B_{k}, C_{k}, D_{k}\right), k=1,2$ be two linear systems. $\alpha_{1}$, $a_{2}$ are said to be similar if there exists a linear continuously invertible operator $W: X_{1} \rightarrow X_{2}$ such that

$$
\begin{aligned}
& A_{2}=W A_{1} W^{-1} \\
& B_{2}=W B_{1} \\
& C_{2}=C_{1} W^{-1} \\
& D_{2}=D_{1}
\end{aligned}
$$

If, moreover, the operator $W$ is unitary, then the systems 1 and 2 are said to be unitarily equivalent.

Definition 2. Let two linear systems $\alpha_{k}=\left(X_{k}, U, V, A_{k}, B_{k}, C_{k}, D_{k}\right), k=1,2$, be such that $U_{2}=V_{1}$.

The linear system $\alpha=(X, U, V, A, B, C, D)$ is called a cascade coupling of $\alpha_{1}, \alpha_{2}$ and is written as $\alpha=\alpha_{2} \alpha_{1}$ if:

$$
\begin{aligned}
& U=V_{1}, V=V_{2}, X=X_{1} \oplus X_{2} \\
& A=A_{1} P_{1}+A_{2} P_{2}+B_{3} C_{1} P_{1} \\
& B=B_{1}+B_{2} D_{1}
\end{aligned}
$$

$$
\begin{aligned}
& C=C_{2} P_{2}+D_{2} C_{1} P_{1} \\
& D=D_{2} D_{1}
\end{aligned}
$$

where $P_{k}$ is the orthoprojection from $X$ onto $X_{k}(k=1,2)$.
We have the following result. If $\alpha=\alpha_{2} \alpha_{1}$, then $\theta_{\alpha}(z)=\theta_{\alpha_{2}}(z) \theta_{\alpha_{1}}(z)$.
Definition 3. The linear system is called an unitary system if the operator

$$
T=\left(\begin{array}{ll}
A & B \\
C & D
\end{array}\right): X \oplus U \rightarrow X \oplus V
$$

is unitary.
We have another equivalent definition.
Definition 4. The linear system is said to be unitary if the following equalities hold:

$$
\begin{align*}
I-A^{*} A & =C^{*} C  \tag{1.1}\\
I-A A^{*} & =B B^{*}  \tag{1.2}\\
I-D D^{*} & =C C^{*}  \tag{1.3}\\
I-D^{*} D & =B^{*} B  \tag{1.4}\\
-A^{*} B & =C^{*} D \tag{1.5}
\end{align*}
$$

An unitary system is also called an unitary colligation [6] or a conservative scattering system [1, 2].

It is known that if $\alpha_{1}$ and $\alpha_{2}$ are unitary systems, then $\alpha_{2} \alpha_{1}$ is also unitary, and according to the Livsis-Brodskii theorem, two simple unitary systems having the same transfer function are unitarily equivalent.

Let us consider the following function model of Nagy and Foias for a simple unitary system constructed by the given transfer function $\theta(z) \in \mathscr{A}(U, V)$.

$$
\begin{aligned}
X & =\left[L_{2}^{+}(V) \oplus \overline{\Delta L_{2}(U)}\right] \ominus\left\{(\theta \omega, \Delta \omega) \mid \omega \in L_{2}^{+}(U)\right\} \\
A(\varphi \oplus \psi) & =e^{-i t}\left(\varphi\left(e^{i t}\right)-\varphi(0)\right) \oplus e^{-i t} \psi\left(e^{i t}\right), \\
B u & =e^{-i t}\left(\theta\left(e^{i t}\right)-\theta(0)\right) u \oplus e^{-i t} \Delta\left(e^{i t}\right) u, \\
C(\varphi \oplus \psi) & =\varphi(0), \\
D u & =\theta(0) u
\end{aligned}
$$

where $\Delta\left(e^{i t}\right)=\left(I-\theta^{*}\left(e^{i t}\right) \theta\left(e^{i t}\right)\right)^{1 / 2} ; \mathscr{A}(U, V)$ denotes the class of all analytic functions in the unit disk $\{z:|z|<1\}$, having values as contractive operators from $U$ to $V$; $L_{2}^{+}(U)$ stands for the Hardy space of elements $f \in L_{2}(U)$ whose $k$ th Fourier coefficient $\hat{f}(k)=0$ for all $k<0$.

Definition 5. Let $\theta(z)=\theta_{2}(z) \theta_{1}(z) ; \theta(z) \in \mathscr{A}(U, V) ; \theta_{k}(z) \in \mathscr{A}\left(U_{k}, V_{k}\right) ; k=1,2$, $U=U_{1} ; V_{1}=U_{2}, V_{2}=V$.

The factorization $\theta(z)=\theta_{2}(z) \theta_{1}(z)$ is said to be regular if

$$
\left\{\overline{\Delta_{2} \theta_{1} h \oplus \Delta_{1} h: h \in L_{2}(U)}\right\}=\overline{\Delta_{2} L_{2}\left(U_{2}\right)} \oplus \overline{\Delta_{1} L_{2}\left(U_{1}\right)}
$$

This definition is equivalent to the following: the operator $\Delta h \mapsto \Delta_{2} \theta_{1} h \oplus$ $\Delta_{1} h$ can be continuously extended to a unitary operator from $\overline{\Delta L_{2}(U)}$ onto $\overline{\Delta_{2} L_{2}\left(U_{2}\right)} \oplus \overline{\Delta_{1} L_{2}\left(U_{1}\right)}$.

From Definition 5, we have the following theorem.
Theorem 1. [12] Suppose the factorization $\theta(z)=\theta_{2}(z) \theta_{1}(z)$ is regular. Then the space

$$
\hat{X}=\left[L_{2}^{+}(V) \oplus \overline{\Delta_{2} L_{2}\left(U_{2}\right)} \oplus \overline{\Delta_{1} L_{2}\left(U_{1}\right)}\right] \ominus\left\{\left(\theta \omega, \Delta_{2} \theta_{1} \omega, \Delta_{1} \omega\right) \mid \omega \in L_{2}^{+}(U)\right\}
$$

contains the subspace

$$
\left.\hat{X}_{2}=\left[L_{2}^{+}(V) \oplus \overline{\Delta_{2} L_{2}\left(U_{2}\right)} \oplus\{0\}\right] \oplus\left\{\theta_{2} u, \Delta_{2} u, 0\right) \mid u \in L_{2}^{+}\left(U_{2}\right)\right\}
$$

invariant for the operator $\hat{A}$ and it is the orthogonal complement in $\hat{X}$ of the subspace

$$
\begin{aligned}
\hat{X}_{1} & =\left\{\left(\theta_{2} u, \Delta_{2} u, v\right) \mid u \in L_{2}^{+}\left(U_{2}\right), v \in \overline{\Delta_{1} L_{2}\left(U_{1}\right)}\right\} \\
& \ominus\left\{\left(\theta_{2} \theta_{1} \omega, \Delta_{2} \theta_{1} \omega, \Delta_{1} \omega\right) \mid \omega \in L_{2}^{+}(U)\right\}
\end{aligned}
$$

## 2. Explicit Cascade Decomposition of a Unitary System According to the Regular Factorization of its Transfer Function

Let a simple unitary system $\alpha=(X, U, V, A, B, C, D)$ be such that its transfer function $\theta(z)$ has a regular factorization $\theta(z)=\theta_{2}(z) \theta_{1}(z)$. We construct explicitly two simple unitary systems $\alpha_{1}$ and $\alpha_{2}$, whose transfer functions are $\theta_{\alpha_{1}}(z)=\theta_{1}(z)$, $\theta_{\alpha_{2}}(z)=\theta_{2}(z)$ and $\alpha=\alpha_{2} \alpha_{1}$, respectively.

According to the Livsis-Brodskii theorem, two simple unitary systems having the same transfer function are unitarily equivalent, so we can use the model of Nagy-Foias for a simple unitary system and still keep the generality of the problem. Besides, the factorization $\theta(z)=\theta_{2}(z) \theta_{1}(z)$ is regular, so instead of the NagyFoias function model $\alpha=(X, U, V, A, B, C, D)$, we can decompose the function model $\hat{\alpha}=(\hat{X}, U, V, \hat{A}, \hat{B}, \hat{C}, \hat{D})$, where

$$
\begin{align*}
\hat{X}= & {\left[L_{2}^{+}(V) \oplus \overline{\Delta_{2} L_{2}\left(U_{2}\right)} \oplus \overline{\Delta_{1} L_{2}\left(U_{1}\right)}\right] } \\
& \ominus\left\{\left(\theta \omega, \Delta_{2} \theta_{1} \omega, \Delta_{1} \omega\right) \mid \omega \in L_{2}^{+}(U)\right\},  \tag{2.1}\\
\hat{A}(\varphi \oplus \psi \oplus \phi)= & \left(e^{-i t}\left(\varphi\left(e^{i t}\right)-\varphi(0)\right) \oplus e^{-i t} \psi\left(e^{i t}\right) \oplus e^{-i t} \phi\left(e^{i t}\right)\right),  \tag{2.2}\\
\hat{B} u= & \left(\frac{\theta\left(e^{i t}\right)-\theta(0)}{e^{i t}} u \oplus \frac{\Delta_{2}\left(e^{i t}\right) \theta_{1}\left(e^{i t}\right)}{e^{i t}} u \oplus \frac{\Delta_{1}\left(e^{i t}\right)}{e^{i t}} u\right),  \tag{2.3}\\
\hat{C}(\varphi \oplus \psi \oplus \phi)= & \varphi(0),  \tag{2.4}\\
\hat{D} u= & \theta(0) u . \tag{2.5}
\end{align*}
$$

By Theorem 1, the subspace $\hat{X}_{2}$ (of $\hat{X}$ ) is invariant for $\hat{A}$ and it is the orthogonal complement in $\hat{X}$ of $\hat{X}_{1}$.

For the spaces $\hat{X}_{1}$ and $\hat{\hat{X}}_{2}$, we construct two systems $\hat{\alpha}_{1}=\left(\hat{X}_{1}, U_{1}, V_{1}, \hat{A}_{1}, \hat{B}_{1}\right.$, $\left.\hat{C}_{1}, \hat{D}_{1}\right)$ and $\hat{\alpha}_{2}=\left(\hat{X}_{2}, U_{2}, V_{2}, \hat{A}_{2}, \hat{B}_{2}, \hat{C}_{2}, \hat{D}_{2}\right)$ as follows:

$$
\begin{align*}
\hat{X}_{1}= & \left\{\left(\theta_{2} u \oplus \Delta_{2} u \oplus v\right) \mid u \in L_{2}^{+}\left(U_{2}\right), v \in \overline{\Delta_{1} L_{2}\left(U_{1}\right)}\right\}  \tag{2.6}\\
& \ominus\left\{\left(\theta_{2} \theta_{1} \omega \oplus \Delta_{2} \theta_{1} \omega \oplus \Delta_{1} \omega\right) \mid \omega \in L_{2}^{+}(U)\right\}
\end{align*}
$$

$\hat{A}_{1}\left(\theta_{2} u \oplus \Delta_{2} u \oplus \Phi\right)=\left(\theta_{2}\left(e^{i t}\right)\left[\frac{u\left(e^{i t}\right)-u(0)}{e^{i t}}\right] \oplus \Delta_{2}\left(e^{i t}\right)\left[\frac{u\left(e^{i t}\right)-u(0)}{e^{i t}}\right] \oplus \frac{\Phi\left(e^{i t}\right)}{e^{i t}}\right)$,

$$
\begin{align*}
& \hat{B}_{1} u=\left(\theta_{2}\left(e^{i t}\right)\left(\frac{\theta_{1}\left(e^{i t}\right)-\theta_{1}(0)}{e^{i t}} u\right) \oplus \Delta_{2}\left(e^{i t}\right)\left(\frac{\theta_{1}\left(e^{i t}\right)-\theta_{1}(0)}{e^{i t}} u\right) \oplus \frac{\Delta_{1}\left(e^{i t}\right)}{e^{i t}} u\right),  \tag{2.8}\\
& \hat{C}_{1}\left(\theta_{2} u \oplus \Delta_{2} u \oplus \Phi\right)=u(0)  \tag{2.9}\\
& \hat{D}_{1} u= \theta_{1}(0) u  \tag{2.10}\\
& \hat{X}_{2}=\left[L_{2}^{+}(V) \oplus \overline{\Delta_{2} L_{2}(U)}+\{0\}\right] \\
& \ominus\left\{\left(\theta_{2} u \oplus \Delta_{2} \omega, 0 \mid u \in L_{2}^{+}\left(U_{2}\right)\right\}\right. \tag{2.11}
\end{align*}
$$

$\hat{A_{2}}(f \oplus g \oplus 0)=\left(e^{-i t}\left(f\left(e^{i t}\right)-f(0)\right) \oplus e^{-i t} g\left(e^{i t}\right) \oplus 0\right)$,
$\hat{C}_{2}(f \oplus g \oplus 0)=f(0)$,

$$
\begin{equation*}
\hat{D}_{2} u=\theta_{2}(0) u \tag{2.14}
\end{equation*}
$$

Then we have the following.
Theorem 2. The simple unitary system of the form (2.1)-(2.5) has an explicit decomposition into two simple unitary systems $\hat{\alpha}_{1}$ and $\hat{\alpha}_{2}, \hat{\alpha}=\hat{\alpha}_{2} \hat{\alpha}_{1}$, by the formulas (2.6)-(2.15).

Proof. Theorem 1 leads to $\hat{X}=\hat{X}_{1} \oplus \hat{X}_{2}$.
Now, we have to prove all operators of the systems $\hat{\alpha}_{1}, \hat{\alpha}_{2}$ are correctly determined.
Firstly, $\hat{A}_{1}$ is an operator from $\hat{X}_{1}$ to $\hat{X}_{1}$. Indeed, since $u \in L_{2}^{+}\left(U_{2}\right)$, then we have $\frac{u\left(e^{i t}\right)-u(0)}{e^{i t}} \in L_{2}^{+}\left(U_{2}\right)$. Moreover, note that for $u \in L_{2}^{+}\left(U_{2}\right)$, and $\Phi \in \overline{\Delta_{1} L_{2}\left(U_{1}\right)}$, $\left(\theta_{2} u, \Delta_{2} u, \Phi\right)$ belongs to the space $\hat{X}_{1}$ if and only if $\theta_{1}^{*} u+\Delta_{1} \Phi$ belongs to $L_{2}^{-}\left(U_{1}\right)$. Hence, for $\left(\theta_{2} u, \Delta_{2} u, \Phi\right) \in \hat{X}_{1}$, we have that

$$
\theta_{1}^{*}\left(\frac{u\left(e^{i t}\right)-u(0)}{e^{i t}}\right)+\Delta_{1} \frac{\Phi\left(e^{i t}\right)}{e^{i t}}=\frac{\theta_{1}^{*}\left(e^{i t}\right) u\left(e^{i t}\right)+\Delta_{1}\left(e^{i t}\right) \Phi\left(e^{i t}\right)}{e^{i t}}-\frac{\theta_{1}^{*}\left(e^{i t}\right) u(0)}{e^{i t}}
$$

belongs to $L_{2}^{-}\left(U_{1}\right)$. Therefore, $\hat{A}_{1}\left(\theta_{2} u, \Delta_{2} u, \Phi\right) \in \hat{X}_{1}$.
Secondly, $\hat{B}_{1}$ is an operator from $U_{1}$ to $\hat{X}_{1}$ because we have

$$
\frac{\theta_{1}\left(e^{i t}\right)-\theta_{1}(0)}{e^{i t}} u \in L_{2}^{+}\left(U_{2}\right)
$$

and

$$
\theta_{1}^{*}\left(e^{i t}\right)\left[\frac{\theta_{1}\left(e^{i t}\right)-\theta_{1}(0)}{e^{i t}} u\right]+\Delta_{1}\left(e^{i t}\right) \frac{\Delta_{1}\left(e^{i t}\right)}{e^{i t}} u=\frac{u}{e^{i t}}-\theta_{1}^{*}\left(e^{i t}\right) \frac{\theta_{1}(0)}{e^{i t}} \in L_{2}^{-}\left(U_{1}\right)
$$

The remainer operators are obviously determined.
Now, we prove $\hat{\alpha}=\hat{\alpha}_{2} \hat{\alpha}_{1}$. It is easy to verify that

$$
\begin{aligned}
&\left(\hat{A}_{1} P_{1}\right.\left.+\hat{A}_{2} P_{2}+\hat{B}_{2} \hat{C}_{1} P_{1}\right)(\varphi, \psi, \phi) \\
&= \hat{A}_{1}\left(\theta_{2} u, \Delta_{2} u, \Phi\right)+\hat{A}_{2}\left(\varphi-\theta_{2} u, \psi-\Delta_{2} u, 0\right)+\hat{B}_{2} \hat{C}_{1}\left(\theta_{2} u, \Delta_{2} u, \Phi\right) \\
&= {\left[\theta_{2}\left(e^{i t}\right)\left(\frac{u\left(e^{i t}\right)-u(0)}{e^{i t}} u\right), \Delta_{2}\left(e^{i t}\right)\left(\frac{u\left(e^{i t}\right)-u(0)}{e^{i t}} u\right), \frac{\phi\left(e^{i t}\right)}{e^{i t}} u\right] } \\
&+\left(\frac{\varphi\left(e^{i t}\right)-\theta_{2}\left(e^{i t}\right) u\left(e^{i t}\right)-\varphi(0)+\theta_{2}(0) u(0)}{e^{i t}}, \frac{\psi\left(e^{i t}\right)-\Delta_{2}\left(e^{i t}\right) u\left(e^{i t}\right)}{e^{i t}}, 0\right) \\
&+\left(\frac{\theta_{2}\left(e^{i t}\right)-\theta_{2}(0)}{e^{i t}} u(0), \frac{\Delta_{2}\left(e^{i t}\right)}{e^{i t}} u(0), 0\right) \\
&=\left(\frac{\varphi\left(e^{i t}\right)-\varphi(0)}{e^{i t}} u, \frac{\psi\left(e^{i t}\right)}{e^{i t}}, \frac{\Phi\left(e^{i t}\right)}{e^{i t}}\right)=\hat{A}(\varphi, \psi, \Phi), \\
&= {\left[\hat{B}_{1}+\hat{B}_{2} \hat{D}_{1}\right) u=\hat{B}_{1} u+\hat{B}_{2}\left(\theta_{1}(0) u\right) } \\
&+\left(\frac{\left.\theta_{2}\left(e^{i t}\right)\left(\frac{\theta_{1}\left(e^{i t}\right)-\theta_{1}(0)}{e^{i t}} u\right), \Delta_{2}\left(e^{i t}\right)\left(\frac{\theta_{1}\left(e^{i t}\right)-\theta_{1}(0)}{e^{i t}} u\right), \frac{\Delta_{1}\left(e^{i t}\right)}{e^{i t}} u\right]}{e^{i t}} \theta_{1}(0) u, \frac{\Delta_{2}\left(e^{i t}\right)}{e^{i t}} \theta_{1}(0) u, 0\right) \\
&=\left(\frac{\theta_{2}\left(e^{i t}\right) \theta_{1}\left(e^{i t}\right)-\theta_{2}(0) \theta_{1}(0)}{e^{i t}} u, \frac{\Delta_{2}\left(e^{i t}\right) \theta_{1}\left(e^{i t}\right)}{e^{i t}} u, \frac{\Delta_{1}\left(e^{i t}\right)}{e^{i t}} u\right) \\
&=\left(\frac{\theta\left(e^{i t}\right)-\theta(0)}{e^{i t}} u, \frac{\Delta_{2}\left(e^{i t}\right) \theta_{1}\left(e^{i t}\right)}{e^{i t}} u, \frac{\Delta_{1}\left(e^{i t}\right)}{e^{i t}} u\right)=\hat{B} u, \\
&\left(\hat{C}_{2} P_{2}+\hat{D}_{2} \hat{C}_{1} P_{1}\right)(\varphi, \psi, \phi) \\
&=\hat{C}_{2}\left(\varphi-\theta_{2} u, \psi-\Delta_{2} u, 0\right)+\hat{D}_{2} \hat{C}_{1}\left(\theta_{2} u, \Delta_{2} u, \Phi\right) \\
&=\varphi(0)-\theta_{2}(0) u(0)+\theta_{2}(0) u(0) \\
& \quad=\varphi(0)-\theta_{2}(0) u(0)+\hat{D}_{2}(u(0))=\varphi(0)=\hat{C}(\varphi, \psi, \phi) \\
& \hat{D}_{2} \hat{D}_{1} u=\theta_{2}(0) \theta_{1}(0) u=\theta(0) u=\hat{D} u .
\end{aligned}
$$

Since $\hat{X}=\hat{X}_{1} \oplus \hat{X}_{2}$ and the operators of the systems $\hat{\alpha}_{1}, \hat{\alpha}_{2}, \hat{\alpha}$ satisfy the equalities of cascade coupling, we have $\hat{\alpha}=\hat{\alpha}_{2} \hat{\alpha}_{1}$.

The system $\hat{\alpha}_{2}$ is constructed according to the function model of Nagy-Foias so $\hat{\alpha}_{2}$ is simple, unitary, and has $\theta_{2}(z)$ as the transfer function. To prove that the system $\hat{\alpha}_{1}$ is simple, unitary, and has $\theta_{1}(z)$ as the transfer function, we consider the system $\alpha_{1}=\left(X_{1}, U_{1}, V_{1}, A_{1}, B_{1}, C_{1}, D_{1}\right)$ constructed by the Nagy-Foias model
corresponding to the transfer function $\theta_{1}(z)$. For the state space $X_{1}$ of $\alpha_{1}$,

$$
X_{1}=\left[L_{2}^{+}\left(V_{1}\right) \oplus \overline{\Delta_{1} L_{2}\left(U_{1}\right)}\right] \ominus\left\{\left(\theta_{1} \omega, \Delta_{1} \omega\right) \mid \omega \in L_{2}^{+}\left(U_{1}\right)\right\}
$$

we consider the operator

$$
\begin{aligned}
\Gamma: X_{1} & \rightarrow \hat{X}_{1} \\
(u \oplus v) & \mapsto\left(\theta_{2} u \oplus \Delta_{2} u \oplus v\right)
\end{aligned}
$$

The operator is determined because we have that $(u \oplus v)$ belongs to $X$ if and only if $\theta_{1}^{*} u+\Delta_{1} v$ belongs to $L_{2}^{-}\left(U_{1}\right)$, and the latter happens if and only if $\left(\theta_{2} u \oplus \Delta_{2} u \oplus v\right)$ belongs to $\hat{X}_{1}$.

Obviously, the operator $\Gamma$ is surjective. Moreover, for any two elements $\left(u_{1} \oplus v_{1}\right),\left(u_{2} \oplus v_{2}\right)$ of $X_{1}$, we have

$$
\begin{aligned}
\left\langle\Gamma\left(u_{1} \oplus v_{1}\right), \Gamma\left(u_{2} \oplus v 2\right)\right\rangle_{\hat{X}_{1}} & =\left\langle\theta_{2} u_{1}, \theta_{2} u_{2}\right\rangle+\left\langle\Delta_{2} u_{1}, \Delta_{2} u_{2}\right\rangle+\left\langle v_{1}, v_{2}\right\rangle \\
& =\left\langle\theta_{2}^{*} \theta_{2} u_{1}+\Delta_{2}^{2} u_{1}, u_{2}\right\rangle+\left\langle v_{1}, v_{2}\right\rangle \\
& =\left\langle u_{1}, u_{2}\right\rangle+\left\langle v_{1}, v_{2}\right\rangle \\
& =\left\langle\left(u_{1} \oplus v_{1}\right),\left(u_{2} \oplus v_{2}\right)\right\rangle_{X_{1}} .
\end{aligned}
$$

Thus, $\Gamma$ is unitary. Besides, it is easy to check that

$$
\Gamma A_{1}=\hat{A}_{1} \Gamma, \Gamma B_{1}=\hat{B}_{1}, C_{1}=\hat{C}_{1} \Gamma, D_{1}=\hat{D}_{1}
$$

Hence, the system $\hat{\alpha}_{1}$ is unitary equivalent to $\alpha_{1}$. Then $\hat{\alpha}_{1}$ is simple, unitary, and has $\theta_{1}(z)$ as the transfer function.

This completes the proof of the theorem.

## 3. Explicit Cascade Decomposition of a Unitary System by its Variant Subspace

Consider a unitary system $\alpha=(X, U, V, A, B, C, D)$ having $X_{2}$ as a subspace invariant for $A$. We construct explicitly two unitary systems $\alpha_{1}$ and $\alpha_{2}$, whose state spaces are $X_{1}=X \ominus X_{2}$, respectively, and $X_{2}$ such that $\alpha=\alpha_{2} \alpha_{1}$.

The idea in this section is based on [6].
By hypothesis, since the subspace $X_{2}$ is invariant for $A$, then $T\left(X_{2}\right)$ is included in $X_{2} \oplus V$ where $T$ is defined in Definition 3. Let $R$ be the orthogonal complement of $T\left(X_{2}\right)$ in the space $X_{2} \oplus V$, i.e., $R=\left[X_{2} \oplus V\right] \ominus T\left(X_{2}\right)$. Now we construct two systems $\alpha_{1}=\left(X_{1}, U_{1}, V_{1}, A_{1}, B_{1}, C_{1}, D_{1}\right)$ and $\alpha_{2}=\left(X_{2}, U_{2}, V_{2}, A_{2}, B_{2}, C_{2}, D_{2}\right)$ with $U_{2}=V_{1}=R$, as follows:

$$
\begin{align*}
& U_{1}=U, V_{2}=V, X_{1}=X \ominus X_{2}  \tag{3.1}\\
& A_{1}=\left.P_{1} A\right|_{X_{1}}  \tag{3.2}\\
& B_{1}=P_{1} B  \tag{3.3}\\
& C_{1}=\left.\left.P_{2} A\right|_{X_{1}} \oplus C\right|_{X_{1}}, \tag{3.4}
\end{align*}
$$

$$
\begin{align*}
& D_{1}=P_{2} B \oplus D  \tag{3.5}\\
& A_{2}=\left.A\right|_{X_{2}}  \tag{3.6}\\
& C_{2}=\left.C\right|_{X_{2}}  \tag{3.7}\\
& B_{2}: R \rightarrow X_{2}, B_{2}\left(x_{2} \oplus v\right)=x_{2}  \tag{3.8}\\
& D_{2}: R \rightarrow V, D_{2}\left(x_{2} \oplus v\right)=v . \tag{3.9}
\end{align*}
$$

Then we have the following.
Theorem 3. The unitary system has an explicit decomposition into two unitary systems $\alpha_{1}, \alpha_{2}$ defined by the formulas (3.1)-(3.9), and $\alpha=\alpha_{2} \alpha_{1}$.
Proof. Obviously, the operators $A_{1}, B_{1}, A_{2}, B_{2}, C_{2}, D_{2}$ are determined. To prove $C_{1}$ and $D_{1}$ are determined, first we observe that, for any $x_{2} \in X_{2}, v \in V, x_{2} \oplus v$ belongs to $R$ if and only if $A_{2}^{*} x_{2}+C_{2}^{*} v=0$. Indeed,

$$
\begin{aligned}
x_{2} \oplus v \in R & \Leftrightarrow\left\langle x_{2} \oplus v, u x_{2}^{\prime}\right\rangle_{x_{2} \oplus V}=0, \quad \forall x_{2}^{\prime} \in X_{2} \\
& \Leftrightarrow\left\langle x_{2} \oplus v, A x_{2}^{\prime} \oplus C x_{2}^{\prime}\right\rangle_{X_{2} \oplus V}=0, \forall x_{2}^{\prime} \in X_{2} \\
& \Leftrightarrow\left\langle x_{2}, A_{2} x_{2}^{\prime}\right\rangle_{X_{2}}+\left\langle v, C_{2} x_{2}^{\prime}\right\rangle_{V}=0, \forall x_{2}^{\prime} \in X_{2} \\
& \Leftrightarrow\left\langle A_{2}^{*} x_{2}+C_{2}^{*} v, x_{2}^{\prime}\right\rangle_{X_{2}}=0, \quad \forall x_{2}^{\prime} \in X_{2} \\
& \Leftrightarrow A_{2}^{*} x_{2}+C_{2}^{*} v=0 .
\end{aligned}
$$

With the operator $C_{1}$ defined by $C_{1} x_{1}=P_{2} A\left(x_{1} \oplus 0\right) \oplus C\left(x_{1} \oplus 0\right)$, we have

$$
\begin{aligned}
A_{2}^{*} P_{2} A\left(x_{1} \oplus 0\right)+C_{2}^{*} C\left(x_{1} \oplus 0\right) & =P_{2} A^{*} A\left(x_{1} \oplus 0\right)+P_{2} C^{*} C\left(x_{1} \oplus 0\right) \\
& =P_{2}\left(A^{*} A+C^{*} C\right)\left(x_{1} \oplus 0\right)
\end{aligned}
$$

by virtue of $\alpha$ being unitary, so the equality (1.1) in Definition 4 leads to

$$
A_{2}^{*} P_{2} A\left(x_{1} \oplus 0\right)+C_{2}^{*} C\left(x_{1} \oplus 0\right)=P_{2} I_{X}\left(x_{1} \oplus 0\right)=0
$$

Thus, $C_{1} x_{1} \in R$ and $C_{1}$ is an operator from $X_{1}$ to $R$. With the operator $D_{1}$ defined by $D_{1} u=P_{2} B u \oplus D u$, it results that $D_{1}$ is an operator from $U$ to $R$. Indeed, from relation (1.5) in Definition 4, we can deduce, for every $u \in U$,

$$
\begin{aligned}
& -A^{*} B u=C^{*} D u \\
\Rightarrow & -P_{2} A^{*} B u=P_{2} C^{*} D u \\
\Rightarrow & -A_{2}^{*} P_{2} B u=C_{2}^{*} D u \\
\Rightarrow & A_{2}^{*} P_{2} B u+C_{2}^{*} D u=0 .
\end{aligned}
$$

Hence, $P_{2} B u \oplus D u \in R$.

Now, we prove that the operators of the systems $\alpha, \alpha_{1}, \alpha_{2}$ satisfy the equalities of cascade coupling:

$$
\begin{aligned}
A_{1} P_{1}+A_{2} P_{2}+B_{2} C_{1} P_{1} & =\left.P_{1} A\right|_{X_{1}}+\left.A\right|_{X_{2}}+B_{2}\left[\left.\left.P_{2} A\right|_{X_{1}} \oplus C\right|_{X_{1}}\right] \\
& =\left.P_{1} A\right|_{X_{1}}+\left.A\right|_{X_{2}}+\left.P_{2} A\right|_{X_{1}} \\
& =\left.A\right|_{X_{1}}+\left.A\right|_{X_{2}}=A, \\
B_{1}+B_{2} D_{1} & =P_{1} B+B_{2}\left[P_{2} B \oplus D\right] \\
& =P_{1} B+P_{2} B=B, \\
C_{2} P_{2}+D_{2} C_{1} P_{1} & =\left.C\right|_{X_{2}}+D_{2}\left[\left.\left.P_{2} A\right|_{X_{1}} \oplus C\right|_{X_{1}}\right] \\
& =\left.C\right|_{X_{2}}+\left.C\right|_{X_{1}}=C, \\
D_{2} D_{1} & =D_{2}\left[P_{2} B \oplus D\right]=D .
\end{aligned}
$$

Finally, we prove that the systems $\alpha_{1}, \alpha_{2}$ defined above are unitary. Consider the system $\alpha_{2}$ with the operator

$$
T_{2}=\left(\begin{array}{ll}
A_{2} & B_{2} \\
C_{2} & D_{2}
\end{array}\right): X_{2} \oplus R \rightarrow X_{2} \oplus V=T\left(X_{2}\right) \oplus R
$$

For every $x_{2} \in X_{2}$, we have

$$
\begin{equation*}
T_{2}\left(x_{2} \oplus 0\right)=A\left(x_{2} \oplus 0\right) \oplus C\left(x_{2} \oplus 0\right)=T_{2}\left(x_{2} \oplus 0\right) \tag{3.10}
\end{equation*}
$$

and, for every $r=x_{2} \oplus v$ belonging to $R$, we have

$$
\begin{equation*}
T_{2}(0 \oplus r)=B_{2} r \oplus D_{2} r=x_{2} \oplus v=r . \tag{3.11}
\end{equation*}
$$

From (3.10) and (3.11), it follows that $T_{2}=\left.T\right|_{X_{2}} \oplus I_{R}$. Since $T$ is unitary, then so is $T_{2}$.

To show that $\alpha_{1}$ is unitary, we first prove

$$
T=\left(T_{2} \oplus I_{X_{1}}\right)\left(I_{X_{2}} \oplus T_{1}\right)
$$

where

$$
T_{1}=\left(\begin{array}{ll}
A_{1} & B_{1} \\
C_{1} & D_{1}
\end{array}\right)
$$

We have

$$
I_{X_{2}} \oplus T_{1}=\left(\begin{array}{ccc}
I_{X_{2}} & 0 & 0 \\
0 & A_{1} & B_{1} \\
0 & C_{1} & D_{1}
\end{array}\right)
$$

is an operator from $X_{2} \oplus X_{1} \oplus U$ to $X_{2} \oplus X_{1} \oplus R$ and

$$
T_{2} \oplus I_{X_{1}}=\left(\begin{array}{ccc}
A_{2} & 0 & B_{2} \\
0 & I_{X_{1}} & 0 \\
C_{2} & 0 & D_{2}
\end{array}\right)
$$

is an operator from $X_{2} \oplus X_{1} \oplus R$ to $X_{2} \oplus X_{1} \oplus V$,

$$
\begin{aligned}
\left(T_{2} \oplus I_{X_{1}}\right)\left(I_{X_{2}} \oplus T_{1}\right) & =\left(\begin{array}{ccc}
A_{2} & 0 & B_{2} \\
0 & I_{X_{1}} & 0 \\
C_{2} & 0 & D_{2}
\end{array}\right)\left(\begin{array}{ccc}
I_{X_{2}} & 0 & 0 \\
0 & \left.P_{1} A\right|_{X_{1}} & P_{1} B \\
0 & \left.\left.P_{2} A\right|_{X_{1}} \oplus C\right|_{X_{1}} & P_{2} B \oplus D
\end{array}\right) \\
& =\left(\begin{array}{ccc}
A_{2} & B_{2}\left(\left.\left.P_{2} A\right|_{X_{1}} \oplus C\right|_{X_{1}}\right) & B_{2}\left(P_{2} B \oplus D\right) \\
0 & \left.P_{1} A\right|_{X_{1}} & P_{1} B \\
C_{2} & D_{2}\left(\left.\left.P_{2} A\right|_{X_{1}} \oplus C\right|_{X_{1}}\right) & D_{2}\left(P_{2} B \oplus D\right)
\end{array}\right) \\
& =\left(\begin{array}{ccc}
A_{2} & \left.P_{2} A\right|_{X_{1}} & P_{2} B \\
0 & \left.P_{1} A\right|_{X_{1}} & P_{1} B \\
C_{2} & \left.C\right|_{X_{1}} & D
\end{array}\right)
\end{aligned}
$$

Since $P_{1} B+P_{2} B=B, \quad C_{2}+\left.C\right|_{X_{1}}=\left.C\right|_{X_{2}}+\left.C\right|_{X_{1}}=C, \quad A_{2}+\left.P_{2} A\right|_{X_{1}}+\left.P_{1} A\right|_{X_{1}}=$ $\left.A\right|_{X_{2}}+\left.A\right|_{X_{1}}=A$, then the matrix above is nothing else but the operator $T=\left(\begin{array}{ll}A & B \\ C & D\end{array}\right)$ from $X \oplus U$ to $X \oplus V$, where $X=X_{2} \oplus X_{1}$. Thus, $\left(I_{X_{2}} \oplus T_{1}\right)=$ $\left(T_{2} \oplus I_{X_{1}}\right)^{-1} T$. From the fact that $T$ and $T_{2}$ are unitary, we conclude that $T_{1}$ is unitary and the proof is complete.

## 4. The Relation Between Cascade Decompositions of a Unitary System by its Transfer Function and by its Invariant Subspace

Note that in the first case of decomposition, given a system with its transfer function $\theta(z)$ having a regular factorization $\theta(z)=\theta_{2}(z) \theta_{1}(z)$, the intermediate space $U_{2}$ is also given. In the second case of decomposition, given a system $\alpha$ with its subspace $X_{2}$ being invariant for $A$, the intermediate space is not given, and we construct this space as $R=\left[X_{2} \oplus V\right] \ominus T\left(X_{2}\right)$. However, we can prove that these two spaces $U_{2}$ and $R$ are unitarily equivalent.

Theorem 4. Let $\alpha=(X, U, V, A, B, C, D)$ be a simple, unitary system constructed according to the Nagy-Foias model. Assume the transfer function $\theta(z)$ of a has a regular factorization $\theta(z)=\theta_{2}(z) \theta_{1}(z), \theta(z) \in \mathscr{A}(U, V), \theta_{i}(z) \in \mathscr{A}\left(U_{i}, V_{i}\right), i=1,2$, $U=U_{1}, V=V_{2}, U_{2}=V_{1}$, and assume $X_{2}$ is a subspace invariant for $A$. Then the space $U_{2}$ is unitarily equivalent to the space $R=\left[X_{2} \oplus V\right] \ominus T\left(X_{2}\right)$ through an operator defined as follows:

$$
\begin{aligned}
& \Gamma: U_{2} \rightarrow R \\
& \Gamma u=\left(e^{-i t}\left(\theta_{2}\left(e^{i t}\right)-\theta_{2}(0)\right) u, e^{-i t} \Delta_{2}\left(e^{i t}\right) u\right) \oplus \theta_{2}(0) u
\end{aligned}
$$

Proof. First, $\Gamma$ is determined since we have

$$
\begin{aligned}
& A_{2}^{*}\left(\frac{\theta_{2}\left(e^{i t}\right)-\theta_{2}(0)}{e^{i t}} u, \frac{\Delta_{2}\left(e^{i t}\right)}{e^{i t}} u\right)+C_{2}^{*} \theta_{2}(0) u \\
& =\left(\left(\theta_{2}\left(e^{i t}\right)-\theta_{2}(0)\right) u-\theta_{2} \omega, \Delta_{2}\left(e^{i t}\right) u-\Delta_{2} \omega\right) \\
& \quad+\left(\theta_{2}(0) u-\theta_{2}\left(e^{i t}\right) \theta_{2}^{*}(0) \theta_{2}(0) u,-\Delta_{2}\left(e^{i t}\right) \theta_{2}^{*}(0) \theta_{2}(0) u\right)
\end{aligned}
$$

where

$$
\begin{aligned}
\omega & =\frac{1}{2 \pi} \int_{0}^{2 \pi}\left[e^{i t} \theta_{2}^{*}\left(e^{i t}\right)\left(\frac{\theta_{2}\left(e^{i t}\right)-\theta_{2}(0)}{e^{i t}}\right) u+e^{i t} \Delta_{2}\left(e^{i t}\right) \frac{\Delta_{2}\left(e^{i t}\right)}{e^{i t}} u\right] d t \\
& =\frac{1}{2 \pi} \int_{0}^{2 \pi}\left[u-\theta_{2}^{*}\left(e^{i t}\right) \theta_{2}(0) u\right] d t=u-\theta_{2}^{*}(0) \theta_{2}(0) u .
\end{aligned}
$$

From this, it follows that

$$
A_{2}^{*}\left(\frac{\theta_{2}\left(e^{i t}\right)-\theta_{2}(0)}{e^{i t}} u, \frac{\Delta_{2}\left(e^{i t}\right)}{e^{i t}} u\right)+C_{2}^{*} \theta_{2}(0) u=0
$$

Hence, $\Gamma u \in R$.
We show that $\Gamma$ preserves the scalar-product. Indeed

$$
\begin{aligned}
\left\langle\Gamma u_{1}, \Gamma u_{2}\right\rangle= & \left\langle\frac{\theta_{2}\left(e^{i t}\right)-\theta_{2}(0)}{e^{i t}} u_{1}, \frac{\theta_{2}\left(e^{i t}\right)-\theta_{2}(0)}{e^{i t}} u_{2}\right\rangle_{L_{2}(V)} \\
& +\left\langle\frac{\Delta_{2}\left(e^{i t}\right)}{e^{i t}} u_{1}, \frac{\Delta_{2}\left(e^{i t}\right)}{e^{i t}} u_{2}\right\rangle_{L_{2}\left(U_{2}\right)}+\left\langle\theta_{2}(0) u_{1}, \theta_{2}(0) u_{2}\right\rangle_{V} \\
= & -\left\langle\theta_{2}\left(e^{i t}\right) u_{1}, \theta_{2}(0) u_{2}\right\rangle+\left\langle\theta_{2}(0) u_{1}, \theta_{2}(0) u_{2}\right\rangle \\
& -\left\langle\theta_{2}(0) u_{1}, \theta_{2}\left(e^{i t}\right) u_{2}\right\rangle+\left\langle\theta_{2}(0) u_{1}, \theta_{2}(0) u_{2}\right\rangle+\left\langle u_{1}, u_{2}\right\rangle \\
= & \left\langle u_{1}, u_{2}\right\rangle .
\end{aligned}
$$

Next, we prove that $\Gamma$ is surjective.
Let $(\varphi, \psi, v) \in R \oplus \Gamma U_{2}$; we will prove $(\varphi, \psi, v)=0$.
Indeed, for any $(\varphi, \psi, v) \in R$, we have $\langle(\varphi, \psi, v), \Gamma u\rangle=0$ for every $u \in U_{2}$ if and only if the conditions below are satisfied:

$$
\begin{gather*}
\theta_{2}^{*} \varphi+\Delta_{2} \psi \in L_{2}^{-}\left(U_{2}\right), \varphi \in L_{2}^{+}(V), \psi \in \overline{\Delta_{2} L_{2}\left(U_{2}\right)},  \tag{4.1}\\
A_{2}^{*}(\varphi, \psi)+C_{2}^{*} v=0,  \tag{4.2}\\
\left\langle\varphi, \frac{\theta_{2}\left(e^{i t}\right)-\theta_{2}(0)}{e^{i t}} u\right\rangle+\left\langle\psi, \frac{\Delta_{2}\left(e^{i t}\right)}{e^{i t}} u\right\rangle+\left\langle v, \theta_{2}(0) u\right\rangle=0 . \tag{4.3}
\end{gather*}
$$

Condition (4.3) is equivalent to

$$
\left\langle e^{i t}\left(\theta_{2}^{*} \varphi+\Delta_{2} \psi, u\right\rangle-\left\langle e^{i t} \varphi, \theta_{2}(0) u\right\rangle+\left\langle v, \theta_{2}(0) u\right\rangle=0\right.
$$

This implies

$$
\langle\omega, u\rangle-\left\langle e^{i t} \varphi, \theta_{2}(0) u\right\rangle+\left\langle v, \theta_{2}(0) u\right\rangle=0
$$

where

$$
\omega=\frac{1}{2 \pi} \int_{0}^{2 \pi} e^{i t}\left(\theta_{2}^{*} \varphi+\Delta_{2} \psi\right) d t \in U_{2}
$$

On the other hand, we have $\varphi \in L_{2}^{+}(V), \theta_{2}(0) u \in U_{2}$. Hence

$$
\left\langle e^{i t} \varphi, \theta_{2}(0) u\right\rangle=0
$$

Then condition (4.3) is equivalent to the condition $\left\langle\omega+\theta_{2}^{*}(0) v, u\right\rangle=0$ for every $u$ belonging to $U_{2}$. This implies $\omega+\theta_{2}^{*}(0) v=0$.

Condition (4.2) is equivalent to

$$
\begin{gather*}
e^{i t} \varphi\left(e^{i t}\right)-\theta_{2}\left(e^{i t}\right) \omega+v-\theta_{2}\left(e^{i t}\right) \theta_{2}^{*}(0) v=0  \tag{4.4}\\
e^{i t} \psi\left(e^{i t}\right)-\Delta_{2}\left(e^{i t}\right) \omega-\Delta_{2}\left(e^{i t}\right) \theta_{2}^{*}(0) v=0 \tag{4.5}
\end{gather*}
$$

From (4.4), it follows that

$$
e^{i t} \varphi\left(e^{i t}\right)+v=\theta_{2}\left(e^{i t}\right)\left[\omega+\theta_{2}^{*}(0) v\right]=0
$$

Therefore, $\varphi\left(e^{i t}\right)=e^{-i t} v \in L_{2}^{-}(V)$. Since $\varphi \in L_{2}^{+}(V)$, we must have $\varphi=0$ and hence, $v=0$.

From (4.5), it follows that

$$
e^{i t} \psi=\Delta_{2}\left[\omega+\theta_{2}^{*}(0) v\right]=0
$$

Hence, $\psi=0$.
We have already proved that $R \ominus \Gamma U_{2}=\{0\}$, so we conclude $\Gamma$ is surjective. Concretely, we can determine

$$
\Gamma^{-1}(\varphi, \psi, v)=\omega+\theta_{2}^{*}(0) v, \text { where } \omega=\frac{1}{2 \pi} \int_{0}^{2 \pi} e^{i t}\left(\theta_{2}^{*} \varphi+\Delta_{2} \psi\right) d t
$$

Thus, $\Gamma$ is unitary and this completes the proof.
Besides, $U_{2}$ and $R$ are unitarily equivalent. We also observe that the operators $A_{2}, B_{2}, C_{2}, D_{2}$ can be considered as the same in both cases of decompositions. Indeed, since $\alpha=\alpha_{2} \alpha_{1}$, then $A_{2}=\left.A\right|_{X_{2}}, C_{2}=\left.C\right|_{X_{2}}$. In the first case of decomposition, we have

$$
B_{2} u=P_{X_{2}}(\Gamma u), D_{2} u=P_{V_{2}}(\Gamma u), B_{2} u+D_{2} u=\Gamma u
$$

while in the second case, we have

$$
B_{2} u=P_{X_{2}} u, D_{2} u=P_{V_{2}} u, B_{2} u+D_{2} u=u
$$

Once $\alpha_{2}$ is constructed, $\alpha_{1}$ is determined uniquely corresponding to $\alpha_{2}$ from the equality $\alpha=\alpha_{2} \alpha_{1}$.

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